Thermalization of electron-boson systems described by a pure state

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Keywords in recent talks

Quantum

Nonequilibrium



Ultrafast experiments in condensed matter

Review: Giannetti et al, Adv. Phys. (2016)

$$|\psi(t)\rangle = e^{-i\hat{H}t}|\psi_0\rangle$$

Two nontrivial outcomes of unitary time evolution





Generate novel states?

States that do not exist in equilibrium phase diagram

Pathway to ergodicity

$$\langle \hat{O}(t) \rangle = \text{Tr}\{\hat{\rho}_{\text{stat}}\hat{O}\}$$

These two concepts are fundamentally different

It underlines the need to understand when and how fast does a condensed-matter system thermalize

Why is understanding of thermalization important in condensed matter?

- Common belief: undriven systems, at asymptotic times after perturbation, should approach a thermal state
- O However, time-resolved experiments may now study the response at extremely short times after perturbation

(in particular ultrafast optics ...)

In a **short time interval**, the dynamics may be efficiently described by models taking into account only the most important interactions

This implies that in a given time interval, the system behaves as a **closed quantum system**

$$|\psi(t)\rangle = e^{-i\hat{H}t}|\psi_0\rangle$$

Closed quantum systems are peculiar:

Many of their properties are implicitly assumed, but rarely verified

Challenges for "Next generation" (1)





Nobel prize 2001

Experimental realization of Bose-Einstein condensation (1995)

"Ultracold atoms in perfectly isolated environment at temperature 20 nK"

Challenges for "Next generation" (1)





Nobel prize 2001

Experimental realization of Bose-Einstein condensation (1995)

"Ultracold atoms in perfectly isolated environment at temperature 20 nK"

Challenges for "Next generation" (2)



"A closed quantum system can never thermalize"

"It is described by a pure state, hence its entropy is always zero"

Closed quantum systems do thermalize

• The notion of **temperature** is valid for generic quantum systems

Review: d'Alessio, Kafri, Polkovnikov, Rigol, Adv. Phys. (2016); and many others

The total entropy of a pure state remains zero forever, however, the entanglement entropy between subsystems increases and reach thermal predictions

Verified experimentally with ultracold bosons

Kaufman et al, Science (2016)



How about condensed-matter systems in pump-probe experiments?

Two extreme views:

The pump pulse creates a thermal electronic distribution, at elevated temperature

(Not entirely correct. However, thermalization can occur very fast ...)

Electrons reach a thermal distribution only after completion of the whole hierarchy of relaxation processes"

(Probably too conservative ...)

Motivation: ultrafast optics

State-of-the-art: Width of the pump pulse ~ 15 fs, Broadband probe at a delay of ~40fs



Modeling of the data consistent with ultrafast relaxation of charge carriers with strongly-coupled excitations of bosonic origin

Thermalization of charrge carriers strongly coupled to a single branch of bosonic excitations ("local" bosons) occurred within 40 fs? Case studied in the following:

Strongly-coupled boson = **dispersionless phonon**

Kogoj, Vidmar, Mierzejewski, Trugman, Bonča, PRB (2016)

Holstein model (single electron)



$$\hat{H} = -t_0 \sum_{j} (\hat{c}_j^{\dagger} \hat{c}_{j+1} + \text{h.c.}) + g \sum_{j} \hat{c}_j^{\dagger} \hat{c}_j (\hat{b}_j + \hat{b}_j^{\dagger}) + \omega_0 \sum_{j} \hat{b}_j^{\dagger} \hat{b}_j$$

Can phonons act as a reservoir/bath?

- Does their spectrum from a continuum?
- Is their intrinsic time scale much shorter than the typical electron time scale?

The system nevertheless does thermalize

"Thinking about thermalization in terms of system + bath is old fashioned"

- Electron and phonons form a closed quantum system
- Simply solve the time-dependent Schroedinger equation exactly

$$|\psi(t)\rangle = e^{-i\hat{H}t}|\psi_0\rangle$$

$$\langle \psi(t) | \hat{O} | \psi(t) \rangle = \text{Tr}\{\hat{\rho}_{\text{stat}} \hat{O}\}$$

Initial state and numerical method



• Unitary time evolution $|\psi(t)
angle=e^{-i\hat{H}t}|\psi_0
angle$

$$\hat{H} = -t_0 \sum_j (\hat{c}_j^{\dagger} \hat{c}_{j+1} + \text{h.c.}) + g \sum_j \hat{c}_j^{\dagger} \hat{c}_j (\hat{b}_j + \hat{b}_j^{\dagger}) + \omega_0 \sum_j \hat{b}_j^{\dagger} \hat{b}_j$$

Ground-state properties

Bonča, Trugman, Batistić, PRB (1999)

Nonequilibrium dynamics

- Vidmar, Bonča, Mierzejewski, Prelovšek, Trugman, PRB (2011)
- Finite-temperature equilibrium

Kogoj, Vidmar, Mierzejewski, Trugman, Bonča, PRB (2016)

Independence of initial state



However:

- Independence of initial state is only a necessary condition
- Demonstrated only for one observable

One-particle density matrix

Goal: to make a statement about all static one-particle correlations



- Momentum distribution function $\hat{n}_k = \frac{1}{L} \sum_{j,l} e^{-i(j-l)k} \hat{c}_j^{\dagger} \hat{c}_l$
- Eigenvalues of the fermionic one-particle density matrix

Thermalization of static electronic correlations on the entire lattice

Dynamic correlations

Result on static observables does not immediately extend to dynamic observables

 $\langle \hat{\jmath}(t)\hat{\jmath}(0)\rangle$

Calculate optical conductivity at time *t* after the quench without applying time-translation invariance



Dynamic correlations

Test thermalization without explicitly carrying out calculations in the Gibbs ensemble

• Thermal equilibrium

$$\sigma_{\rm reg}'(\omega) = \frac{1 - e^{-\omega/T}}{\omega} C(\omega)$$

$$C(\omega) = \Re \int_0^\infty \mathrm{d}t e^{i\omega^+ t} \mathrm{Tr}\{\hat{\rho}_{\rm sys}\hat{\jmath}(t)\hat{\jmath}(0)\}$$

$$R(\omega) = \frac{C(\omega) - C(-\omega)}{C(\omega) + C(-\omega)} = \tanh\left(\frac{\omega}{2T}\right)$$

$$C(\omega,t) = \Re \int_0^\infty \mathrm{d}s e^{i\omega^+ s} \langle \psi_0 | \hat{j}(t+s)\hat{j}(t) | \psi_0 \rangle$$

$$R(\omega, t) = \frac{C(\omega, t) - C(-\omega, t)}{C(\omega, t) + C(-\omega, t)}$$



Dynamic correlations

Test thermalization without explicitly carrying out calculations in the Gibbs ensemble



Static and dynamic correlations - Temperature

Static fermionic correlations: temperature obtained from the Gibbs ensemble by matching the electron kinetic energy

Dynamic correlations: temperature obtained by fitting $\langle R(\omega,t) \rangle_t$ with the thermal form $\tanh(\omega/(2T))$

Are these temperatures equal? Yes!



The temperature, measured in response functions, is the temperature of the closed electron-phonon system



Simple, closed quantum systems may thermalize extremely fast

(Useful input for ultrafast optical experiments)

• "Thinking about thermalization in terms of system + bath is old fashioned"

(Is it really true? Find more examples ...)

Thank you!