

in the spin-orbital systems

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Acknowledgments

Orbital excitons / cuprates:

Justina Schlappa, Cheng-Chien Chen, Maria Daghofer, Maurits Haverkort, Thorsten Schmitt,

Valentina Bisogni, Jeroen van den Brink, et al.

Spin-orbital excitons / iridates:

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Putting 1 hole into the 1D antiferromagnet (AF, ground state of the undoped 1D Hubbard model):



- \rightarrow hole (~holon) + domain wall (~spinon) separate
- → paradigm: spin-charge separation in 1D [T. Giamarchi, Quantum Physics in One Dimension (2004)]

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→ observed by ARPES on undoped quasi-1D cuprates [C. Kim et al., PRL 77, 4054 (1996)]

Putting 1 hole into the 2D AF (ground state of the undoped 2D Hubbard model):



 \rightarrow hole (~holon) excites collective magnetic excitations (~magnons) when moving

 \rightarrow not only spin and charge does *not* separate but even... holon motion hindered by magnons

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Outline

1. Spin-orbital separation in quasi-1D cuprates:

- theory
- experiment
- postsciptum (PS)
- 2. Polaronic motion of j=3/2 spin-orbital excitons in quasi-2D iridates:
 - theory
 - experiment
 - postscriptum (PS)
- 3. Conclusions

Single Cu^{2+} ion in Sr₂CuO₃ (1 hole in 3*d* orbitals):

crystal field \rightarrow hole with s=1/2 spin in the x^2-y^2 orbital (ground state) & 4 excited orbitals



1D lattice of Sr₂CuO₃:

hopping + Coulomb repulsion \rightarrow low energy: Heisenberg superexchange between s=1/2 spins

$$\mathcal{H} = J \sum_{\langle ij \rangle} \left(\mathbf{S}_i \mathbf{S}_j + \frac{1}{4} \right)$$

[KW et al., PRL 107, 147201 (2011); KW et al., PRB 88, 195138 (2013)]

How does orbital excitation move in 1D s=1/2 AF?

1st step: 1D AF and ferroorbital (FO) ground state



[KW et al., PRL 107, 147201 (2011); KW et al., PRB 88, 195138 (2013)]

How does orbital excitation move in 1D s=1/2 AF?

1st step: 1D AF and FO

2nd step: we create orbital excitation (also called: exciton)



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1st step: 1D AF and FO

2nd step: we create orbital excitation

3rd step: orbital excitation moves and creates 1 spinon



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Spin of electron in the upper orbital is *conserved*

during this superexchange process

(a rather realistic assumption):





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4th step: further motion does *not* create more spinons



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spin-orbital separation



How does orbital excitation move in 1D s=1/2 AF?



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1. Spin-orbital separation in quasi-1D cuprates: experiment

Resonant inelastic x-ray scattering (RIXS) at Cu L_3 edge in Sr₂CuO₃





[J. Schlappa et al., Nature 485, 82 (2012); V. Bisogni et al., PRL 114, 096402 (2015)]

1. Spin-orbital separation in quasi-1D cuprates: experiment

Excellent agreement with the experiment and theory (~spin-orbital model)



[J. Schlappa et al., Nature 485, 82 (2012); V. Bisogni et al., PRL 114, 096402 (2015)]

1. Spin-orbital separation in quasi-1D cuprates: experiment

But where is the 'pure' orbiton?



[J. Schlappa et al., Nature 485, 82 (2012); V. Bisogni et al., PRL 114, 096402 (2015)]

The above physics is only valid when strong crystal field fully polarizes the ground state (FO state)



Single Ir^{4+} ion in $Sr_{1}IrO_{4}$ (1 hole in 5*d* orbitals):

crystal field + spin-orbit \rightarrow hole in *j*=1/2 spin-orbital isospin ground state and *j*=3/2 excitations



2D lattice of Sr₂IrO₄:

hopping + Coulomb repulsion \rightarrow low energy: Heisenberg superexchange between j=1/2 isospins

$$\mathcal{H} = J \sum_{\langle ij \rangle} \left(\mathbf{S}_i \mathbf{S}_j + \frac{1}{4} \right)$$

 1^{st} : ground state is a 2D AF formed by j=1/2 isospins



2. Polaronic motion of j=3/2 spin-orbital excitons in quasi-2D iridates: theory

How does j=3/2 spin-orbital excitation move in 2D j=1/2 AF?

 2^{nd} : we create a single j=3/2 excitation in the ground state



 3^{rd} : we propagate it to the *nn* site via superexchange process $\rightarrow a j=1/2$ magnon left behind



 3^{rd} : we propagate it to the *nn* site via superexchange process $\rightarrow a j=1/2$ magnon left behind



 4^{th} : we propagate the excitation further \rightarrow *more* magnons left behind



 4^{th} : we propagate the excitation further \rightarrow *more* magnons left behind



Only motion by coupling to spin fluctuations possible, just like for a hole in 2D AF

 \rightarrow also in this case the *j*=3/2 exciton moves as a polaron

2. Polaronic motion of j=3/2 spin-orbital excitons in quasi-2D iridates: experiment

Good agreement between experiment and theory

Ir
$$L_3$$
 edge RIXS on Sr₂IrO₄

Self-consistent Born approximation calculations



Note: theoretical calculations include the polaronic motion of the j=3/2 excitons

2. Polaronic motion of j=3/2 spin-orbital excitons in quasi-2D iridates: PS

Is that the full story?





What is the origin of the small branch of the exciton dispersion with the minimum at Γ ?

2. Polaronic motion of j=3/2 spin-orbital excitons in quasi-2D iridates: PS

Not in the previous model: the Jahn-Teller interaction between j=3/2 and j=1/2 isospins



[E. M. Plotnikova et al.,, Phys. Rev. Lett. 116, 106401 (2016)]

2. Polaronic motion of j=3/2 spin-orbital excitons in quasi-2D iridates: PS

This allows for the hopping of j=3/2 exciton which does not introduce defects in AF...



2. Polaronic motion of j=3/2 spin-orbital excitons in quasi-2D iridates: PS

... and it may explain the extra feature with minimum at Γ in RIXS (peak "A")



3. Conclusions

1. Motion of orbital exciton in quasi-1D cuprates:

spin-orbital separation, just like spin-charge separation (note: valid only for strong crystal field)

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2. Motion of spin-orbital exciton in quasi-2D iridates:

polaronic type of motion, just like for a hole in 2D AF (though Jahn-Teller changes it a bit)



Take-home message

Systems *without* strong on-site spin-orbit coupling ("3d": cuprates, manganites, etc.)

are NOT always very different from

the ones with strong on-site spin-orbit coupling ("5d": iridates, osmates, etc.)



[W. Witczak-Krempa, Annual Review of Condensed Matter Physics, 5, 57 (2014)]