



CBB Workshop at UoC, 27-28/10 2017, Chicago, IL

## **Benchmarking Coherent** Synchrotron Radiation

### S. Di Mitri, ELETTRA SINCROTRONE TRIESTE

### INNOVATIONS IN BRIGHT BEAM SCIENCE





# experimental evidences.

## spread, microbunching). New paths of research on CSR instability could be identified.

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## Elettra Sincrotrone Trieste

- I will skip theoretical derivations and computational algorithms.
- □ I will report about a selection of case studies for **linac-driven FELs**.

### □ This is a review with an accent on **analytical modelling**, and **accuracy of predictions** relative to

### $\Box$ CSR will be analysed in terms of transverse ( $\rightarrow$ emittance) and longitudinal instability ( $\rightarrow$ energy)





### Two particles on same circular path (1-D model), • $\gamma^3 s/R \gg 1$ (steady-state approximation),

р

### ENERGY CHANGE ALONG BUNCH, per METER:



## Current spikes or fast rises enhance the z-CSR field.

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## **CSR Tail-Head Interaction**

P receives photons (test particle in the bunch head)

Curved path in a dipole magnet

## **O** The longitudinal E-field comes from geometrical considerations.

## effect.

### **D** Effect on transverse emittance can be described through linear optics.



**CSR** instability is primarily a **chromatic** 





### Arc (TBA or Bates cell) 0.1 GeV, 0.1 nC, < 1 kA

## Energy Loss

- Example: Gaussian bunch in steady-state emission:  $U_{tot} = -0.028 \times e^2 Z_0 c N \frac{\theta R^{1/3}}{\sigma_a^{4/3}} = -0.16 \text{ MeV} \text{ Total energy loss}$



### C. Hall et al., PRST-AB 18, 030706 (2015)

 $\sigma_z = 50 \mu m$ Q = 300 p C $L_B = 1 m$  $\theta_B = 10^{o}$ R = 5.7m $I_{pk} = 715 A$ E = 700 MeV





### • Particle coordinates transform according to:





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## **Collapsing the Physics onto RMS Values**

## Change of longitudinal momentum



Is this simplified picture good enough? At which level of accuracy?

### RMS EMITTANCE

 $\mathcal{E}_{x} = \sqrt{\left\langle x_{\beta}^{2} \right\rangle \left\langle x_{\beta}^{2} \right\rangle - \left\langle x_{\beta} x_{\beta}^{\prime} \right\rangle^{2}}$ 



 $\varepsilon_x^2 = \varepsilon_{x,0}^2 + \varepsilon_{x,0} \left( \beta_x \left\langle \Delta x'^2 \right\rangle + 2\alpha_x \left\langle \Delta x \Delta x' \right\rangle + \gamma_x \left\langle \Delta x^2 \right\rangle \right) + \varepsilon_{x,0} \left( \beta_x \left\langle \Delta x'^2 \right\rangle + 2\alpha_x \left\langle \Delta x \Delta x' \right\rangle + \gamma_x \left\langle \Delta x^2 \right\rangle \right) + \varepsilon_{x,0} \left( \beta_x \left\langle \Delta x'^2 \right\rangle + 2\alpha_x \left\langle \Delta x \Delta x' \right\rangle + \gamma_x \left\langle \Delta x^2 \right\rangle \right) + \varepsilon_{x,0} \left( \beta_x \left\langle \Delta x'^2 \right\rangle + 2\alpha_x \left\langle \Delta x \Delta x' \right\rangle + \gamma_x \left\langle \Delta x^2 \right\rangle \right) + \varepsilon_{x,0} \left( \beta_x \left\langle \Delta x'^2 \right\rangle + 2\alpha_x \left\langle \Delta x \Delta x' \right\rangle + \gamma_x \left\langle \Delta x^2 \right\rangle \right) + 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### Since $\Delta x$ and $\Delta x'$ from CSR field are correlated, this goes to 0.





## Horizontal Emittance



![](_page_5_Figure_10.jpeg)

### S. Di Mitri et al., PRL 110, 014801 (2013) T. Hara @ FEL'17 and NOCE'17

![](_page_6_Picture_0.jpeg)

## lengths $L_t \simeq (24R^2\sigma_z)^{1/3} \simeq 0.1 - 1 \text{ m}$

![](_page_6_Picture_2.jpeg)

![](_page_6_Figure_3.jpeg)

## **Transient Effects**

![](_page_6_Figure_7.jpeg)

![](_page_6_Picture_8.jpeg)

![](_page_6_Picture_9.jpeg)

### Steady-state model doesn't account for transient effects, which are relevant over

### Y. Jiao et al., PRST-AB 17, 060701 (2014) S. Di Mitri, NIM A 806 (2016)

![](_page_7_Picture_0.jpeg)

## CSR kicks are calculated

![](_page_7_Figure_3.jpeg)

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## I. Akkermans et al., PRST-AB 20, 080705 (2017) **Example: Compact ERL – UV FEL**

![](_page_8_Picture_0.jpeg)

![](_page_8_Figure_2.jpeg)

## 1-D Codes

S. Di Mitri et al., NIM A 608 (2009) K. Bane et al., PRST-AB 12, 030704 (2009)

> FERMI Team, P. Williams, A. Brynes & ASML, in progress

1-D codes agree at full compression even (??). 1-D approx. still seems to be good for  $\sigma_1 \approx (R\sigma_z^2)^3$ 

0.3 GeV, 0.7 nC, < 3 kA

![](_page_9_Picture_0.jpeg)

![](_page_9_Figure_2.jpeg)

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## 2,3-D Codes

□ 1-D model doesn't account for CSR field radial dependence: • Forcing  $\theta_1 \cong \theta_2$  leads to the so-called Derbenev criterion: transverse effects become important when  $\sigma_{\perp} \gg (R \sigma_z^2)^{1/3}$ 

3-D model

![](_page_9_Figure_9.jpeg)

![](_page_9_Figure_10.jpeg)

![](_page_9_Figure_11.jpeg)

### H. Owen et al. for NLS (2008) C. Prokop et al., NIM A 719 (2013)

1-D approximation

![](_page_10_Picture_0.jpeg)

- □ Slice emittance growth becomes noticeable in ultra-low emittance beams. 3-D effects are weakened in "pencil" beams... In the second Slice linear optics mismatch Chromatic aberrations

### 0.2 GeV, 0.2 nC, < 0.2

![](_page_10_Figure_8.jpeg)

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## Sincrotrone Trieste

0.2

-ĭ.5

-0.5

s (m)

![](_page_10_Figure_14.jpeg)

0.5

 $x 10^{-4}$ 

### S. Bettoni et al., PRAB 19, 034402 (2016)

![](_page_11_Picture_0.jpeg)

### □ Shielding of CSR field would require pipe gap as small as < 2 mm or so.

### J. Esberg et al. for CERN (2015)

![](_page_11_Figure_5.jpeg)

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## Shielding

![](_page_11_Figure_14.jpeg)

![](_page_11_Figure_17.jpeg)

FIG. 14. (Color) Realistic magnets: Parameter set E (JLab TH2) magnet) line (top), set F (CESR analyzer magnet) (bottom). Bmad agrees with the CSR-wake formula Eq. (53) better than the other codes at the bunch tail.

### V. Yakimenko et al. @ ATF (2012) Model Exp. Plates gap [mm] Plates gap [mm]

Energy spectrum [KeV]

![](_page_12_Picture_0.jpeg)

Bunching is proportional to this parameter when CSR only is considered.

![](_page_12_Picture_2.jpeg)

![](_page_12_Picture_4.jpeg)

## **CSR-Microbunching: Theory**

dipole line, CSR can even dominate the instability gain. □ Theoretical prescriptions on beam optics were found for minimizing both the CSR-induced

![](_page_12_Figure_7.jpeg)

## C.-Y. Tsai et al., PRAB 20, 024401 (2017)

- CSR amplifies microbunching gain typically driven by longitudinal space charge force. In a multi
  - emittance growth and microbunching gain (local isochronicity,  $\pi$ -phase advance, small betas, ecc.).

![](_page_13_Picture_0.jpeg)

- - Picture is unclear to me yet

# D Proposal of research lines:

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## **Final Remarks**

□ 1-D steady-state analytical formulas guide to the design of e-beam lines (chicanes, transfer lines, arcs). •  $\sigma_{\delta} \sim 0.1\%$ ,  $\Delta \varepsilon_{n,x} \sim 0.1 \ \mu m$  accuracy of prediction for E > 300 MeV,  $\sigma_{z} > 10s \ \mu m$ , Q < 0.5 nC (I < kAs) Control of CSR-induced emittance growth through beam optics is well-established Predictions get worse at lower beam energies, and with multiple bends in long beam lines

1-D approximation seems to be ok for a relaxed Derbenev criterion even 3-D effects are expected to raise at full compression

Systematic investigation of 1-D vs. 3-D effects (codes benchmarking, experimental accuracy) Validation of the CSR-driven microbunching gain vs. beam optics Direct characterization of CSR-induced distortions in (z,E), (z,x) and (z,x') phase space.

![](_page_13_Picture_15.jpeg)

![](_page_14_Picture_0.jpeg)

![](_page_14_Picture_1.jpeg)

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Thank you for Your attention